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# Edges, Transitions and Criticality

Suman Kumar Maji\* and Hussein M. Yahia

*Geostat Team (Geometry and Statistics in Acquisition Data) of INRIA,  
200 rue de la Vieille Tour, 33405 Talence Cedex, FRANCE*

*Email: sumankumar.maji, hussein.yahia@inria.fr*

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## Abstract

In this article, various notions of edges encountered in digital image processing are reviewed in terms of compact representation (or completion). We show that critical exponents defined in Statistical Physics lead to a much more coherent definition of edges, consistent across the scales in acquisitions of natural phenomena, such as high resolution natural images or turbulent acquisitions. Edges belong to the multiscale hierarchy of an underlying dynamics, they are understood from a statistical perspective well adapted to fit the case of natural images. Numerical computation methods for the evaluation of critical exponents in the non-ergodic case are recalled, which apply for the vast majority of natural images. We study the framework of reconstructible systems in a microcanonical formulation, show how it redefines edge completion, and how it can be used to evaluate and assess quantitatively the adequation of edges as candidates for compact representations. We study with particular attention the case of turbulent data, in which edges in the classical sense are particularly challenged. Tests are conducted and evaluated on a standard database for natural images. We test the newly introduced compact representation as an ideal candidate for evaluating turbulent cascading properties of complex images, and we show better reconstruction performance than the classical tested methods.

*Keywords:* Edge detection, complex systems, statistical physics, critical exponents, compact representation, nonlinear signal processing, multifractals, wavelets, multiresolution analysis, turbulence, scale space.

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\*Corresponding author. Tel.: (+33)(0)5 24 57 40 49, Fax : (+33)(0)5 24 57 41 24.

## 1. Introduction

As algorithms dedicated to the computation of edges in digital images emerged [1, 3, 4, 7, 5, 6, 9, 10, 11, 12, 13, 14, 36, 15], Torre and Poggio [16], while observing that most methods rely on the ill-posed problem of differentiating digital images, proposed a general qualitative description of *edges*: they noted that edges are naturally associated to the concepts of *compact representation* (they call it *completion*), i.e. edges encode most information of an image [57]. Similarly, other authors note that edges represent an image's independent features [20]. In [16] the authors focus on edge detection as the process of computing derivatives, and, while attempting to do so in a well-posed form, they are led naturally to the problem of prefiltering the image by a (e.g. Gaussian) kernel, which transforms the input signal into a differentiable mapping in the continuous domain, hence allowing the characterization of edges by differential operators. An instance of this formalism is the zero-crossing of second-order derivatives, as in [10, 17, 11, 14, 52, 55, 56], or [44] to cite few, including a recent nonlinear derivative approach (called NLFS) [53]. This formal setting allowed the development of edge characteristics in the framework of differential geometry, a perspective that has become pervasive in image processing [18, 34]. The multiscale nature of edges was recognized very early and it was noted that tracing edge properties across scales would gain insight into the physical process behind image formation. Neurophysics was demonstrating that, in the optical pathway, spatial filters of different sizes operate at the same location [19]. This is related to the processing of information in the early visual system [24], where cells tend to take advantage of the statistical regularities of the input signal in order to get compact representations out of redundancy [25, 26, 27].

The convolution of the input image signal by a Gaussian kernel introduces a scale parameter (the standard deviation of the Gaussian kernel) corresponding to a simple linear scale-space associated to the heat equation. This is often used as an argument for advocating multiscale properties of Gaussian prefiltering [16, 29, 32, 33, 37]. In general, however, the multiscale properties of complex systems do not comply with such an extreme simplification [28]. The advent of scale-space theory in Computer Vision allowed more complex multiscale representations corresponding, among others, to anisotropic diffusion schemes [30, 35, 46, 54], which can incorporate probabilistic models of both sensor noise and operators' responses (to better estimate the gradient's magnitude threshold in case of noise [47]). However the simple example

of an image corresponding to the acquisition of a turbulent fluid, like, for instance, a remotely sensed acquisition over the oceans, contains coherent structures associated to the cascading properties of intensive variables in Fully Developed Turbulence (FDT) [42]. It has an associated multiscale hierarchy consisting of sets having a multifractal nature [31] and, as such, cannot be contemplated within a differentiable scale-space framework. Incidentally note that in [49] authors write that an appropriate spatial scale depends upon the local structure of the edge, and thus varies *unpredictably* over the image.

In a seminal paper, Mallat and Zhong [38] relate multiscale Canny edge detection to the local maxima of a wavelet transform and study the completion of multiscale edges associated to the maxima of wavelet coefficients (multiscale edge detection [38, 58]); for that purpose they introduce a reconstruction algorithm of a signal from its edges. The quality of the reconstruction is quantitatively evaluated by the SNR of the original and reconstructed signals, hence providing an accurate evaluation of the quality of *edge pixels* determined by their method; this evaluation is different than the previous criteria used in computing edge detection performance [18, 43]. Local maxima of wavelet coefficients are also used by other authors to form the basis of the Wavelet Transform Modulus Maxima (WTMM) methodology [40]. Edges can also be understood as alignment of Fourier or wavelet phases across scales [21, 22, 23].

In this paper, we show that recent developments around the notion of *transition* in nonlinear physics, along with enhanced computational methods of its quantitative parameters (most notably *singularity exponents*) [42], lead to a notion of edge that provides better results over all the previous declined versions encountered in image processing w.r.t. edge completion. Our results strongly advocate for a definition of edge based on nonlinear operators while we prove along the way, and incidentally, that previous nonlinear approach [53] also works better than the classical ones w.r.t. to reconstruction of an image from its edge data. This can be done by referring to the early-addressed pertinent notion of *compact representation* (completion). When neurophysics and the study of biological vision in mammals state that edges encode most information in an image signal, this must have the consequence of being able to reconstruct accurately an image from the compact representation of its edge pixels ([58], p. 194). Statistically, image acquisitions correspond to processes out of the equilibrium state, so that, from a theoretical point of view, transitions associated to scale-space formulations reviewed



above cannot be correct. We show that the new notion of edge outperforms the most well-known previous ones and that it is naturally robust to noise. We give specific attention to the case of turbulent images, whose edges are not well defined in the classical context of edge detection, and we show that in this context the new notions introduced in this article work much better than the previous ones.

Unpredictability of edges, the concept of singularity exponents and the framework of reconstructible systems are introduced in section 2. Edge consistency across the scales is addressed in section 3. Quantitative results are shown and discussed in section 4 and section 5, where in subsection 4.4 the case of turbulent images is specifically addressed. Conclusion (section 6) follows.

## 2. Edges, unpredictability and reconstructibility revisited using the microcanonical multiscale formalism (MMF)

In this section we show how the microcanonical multiscale formalism (MMF [71]) can be applied to edge detection and image reconstruction. We will show that ideas borrowed from Statistical Physics about criticality and exponents, when evaluated in a microcanonical formulation, are associated to a computable notion of *transition*, intimately related to *predictability* in complex systems [59]. This in turn gives rise to a notion of edge whose quantitative performance evaluation can be tested through the framework of *reconstructible systems*.

### 2.1. Local Predictability Exponents

Classically, edges are related to sharp variations of the image gradient. The main idea worked out in this article is to delve deeper into developing the notion of “sharp variation”, and relate it to the more general notion of “transition” defined in Statistical Physics for *intensive* physical variables. For that matter, a scalar image  $I$  defined over a compact subset of  $\mathbb{R}^2$  is identified with an intensive physical variable (such an identification corresponds exactly to the physics of the acquisition for images of natural phenomena such as in infrared remote sensing imagery etc.). In nonlinear physics, the relation between the transitions of an intensive variable and criticality is well known, explained, and quantitatively formalized through the notion of *critical exponent* [62]. We recall that definition here. We say that image  $I$  has

a critical exponent  $h(\vec{x})$  at point  $\vec{x}$ , if for at least one multiscale functional  $\mathbb{T}_r$ <sup>1</sup>, dependent on scale  $r$ , the following equation holds:

$$\mathbb{T}_r I(\vec{x}) = \alpha(\vec{x}) r^{h(\vec{x})} + o(r^{h(\vec{x})}) \quad (r \rightarrow 0) \quad (1)$$

where  $\alpha(\vec{x})$  is a signal-dependent amplitude prefactor. An effective choice for the functional  $\mathbb{T}_r$  leads to a measure given by the total variation of the image gradient [71], and is defined as:

$$\mu(\mathcal{B}_r(\vec{x})) = \int_{\mathcal{B}_r(\vec{x})} \|\nabla I\|(\vec{x}) d(\vec{x}) \quad (2)$$

where  $\mathcal{B}_r(\vec{x})$  is a ball of radius  $r$  centered at point  $(\vec{x})$  of the signal domain. The wavelet transform of the measure  $\mu$  then allows us to infer a more computable version of the singularity exponents (less prone to noise variation) [72], such that:

$$\mathcal{T}_\Psi \mu(\vec{x}) = \alpha_\Psi(\vec{x}) r^{h(\vec{x})} + o(r^{h(\vec{x})}) \quad (r \rightarrow 0) \quad (3)$$

where  $\Psi$  is an associated wavelet.

The central problem is to compute at high numerical precision the value of  $h(\vec{x})$  at point  $\vec{x}$ : bad approximations of singularity exponents lead to poor reconstructions. We will address the problem of robust computation of singularity exponents in section 2.2. For the moment we note that equation (1) is a pointwise and localized version of the definition used in introducing singularity spectrum [40, 70]: we don't make use of statistical averages and grand ensembles, but seek to evaluate  $h(\vec{x})$  at point  $\vec{x}$  (which means we cannot rely on stationarity hypothesis). We denote  $\mathcal{F}_h$  the component in the image's domain associated to singularity exponent value  $h$  as:

$$\mathcal{F}_h = \{\vec{x} : h(\vec{x}) = h\} \quad (4)$$

This family of sets is naturally associated to the multiscale hierarchy in an image [71]. In the case of natural images of the physical world, it is expected that the values of  $h(\vec{x})$  are bounded by below, so that the *Most Singular Manifold* or MSM can be defined as:

$$\mathcal{F}_\infty = \{\vec{x} : h(\vec{x}) = h_\infty = \text{Min}(h(\vec{x}))\} \quad (5)$$

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<sup>1</sup>Typical examples of  $\mathbb{T}_r$  can be wavelet transforms at scale  $r$  or some differential operators applied to the signal.

noting that, in digital signals, the value  $h_\infty$  is thresholded and must correspond to a (small) tolerance interval. The MSM comprises the set of points in an image with the sharpest transitions, corresponding to most unpredictable points. Consequently, if this set is computed properly, it must maximize the statistical information content in the image, so that we should be able to reconstruct the signal from its restriction to this set: this is the framework of *reconstructible systems* that will be our fundamental tool in evaluating the quality of a compact representation for a set of edges.

Criticality, and the associated *critical exponents* are key notions in Statistical Physics to understand phase transitions, which are prototypes of scale invariant phenomena. In high order transitions, thermodynamical variables as well as time and spatial correlation functions display a power-law behaviour, where the characteristics of the underlying system are determined completely, at the macroscopic scale, by the value of the exponent [62]. Critical exponents are universal: systems having the same values of critical exponents belong to a same *universality class* and their statistical macroscopic multiscale characteristics are similar, independantly of the causalities in the extended phase space. This is used in recent advances in complex systems and complex network methods, where fingerprints of universality classes are checked from singularity spectra computed from critical exponents [69], or the statistics in natural images are found identical to those of complex systems [24]. The distribution of critical exponents in a system determines its multiscale properties which are accessible statistically. In particular, the classical notion of gradient, which serves as the basic ingredient in most edge detection algorithm is not clearly defined in such context and one must find more accurate statistical description of transitions across the scales.

## 2.2. Computation of the singularity exponents

Following equation (1),  $h(\vec{x})$  can be obtained by a linear regression of  $\log(\mathbb{T}_r I(\vec{x}))$  vs.  $\log(r)$ :

$$h(\vec{x}) = \lim_{r \rightarrow 0} \frac{\log(\mathbb{T}_r I(\vec{x})/\alpha(\vec{x}))}{\log r} \quad (6)$$

A very fast but crude version of computing  $h(\vec{x})$  is known as the Gradient histogram method [41], which takes into account the multiscale measure defined in equation (2), at a minimum resolution  $r_0$ . The scale  $r_0$  is chosen such that the whole image corresponds to size 1; in other words, if the image

is an array of discretize values of size  $m \times n$ , one chooses  $r_0 = 1/\sqrt{m \times n}$ . Approximating  $\alpha(\vec{x})$  as the average of the norm of the gradients, an estimate of  $h(\vec{x})$  is obtained as:

$$\tilde{h}(\vec{x}) \equiv \frac{\log(\|\nabla I\|(\vec{x})/\langle\|\nabla I\|\rangle)}{\log r_0} \quad (7)$$

The method however fails to produce satisfactory results for natural images and is highly sensitive to noise [41]. A better solution would consist in performing the regression on wavelet projection of measures [71], i.e., a logarithmic regression applied to equation (3). However, a wavelet projection of the measurement at various scales is costly in computation time and only serves to enhance the resolution of less singular structures at the cost of coarsening the most singular ones [41]. Since the objective is to recover the most singular structures, a better optimized and cost-effective way is to use a point estimation of the singularity exponents. From the perspective of reconstructible systems, good evaluation algorithms come from the observation that the set of most unpredictable points  $\mathcal{F}_\infty$  (see equation (5)) that provides a perfect reconstruction is such that [72]:

$$\text{div}(\nabla I|_{\mathcal{F}_c}) = 0. \quad (8)$$

where  $\mathcal{F}_c$  is the complementary set of  $\mathcal{F}_\infty$ . As a consequence, singularity exponents can be called *Local Predictability Exponents*, they encode predictability information [59], like Lyapunov exponents, and are better evaluated in digital signals by one of the following approximation:

$$h(\vec{x}) = \frac{\log \mathcal{T}_\Psi \mu(\vec{x}, r_0)/\langle \mathcal{T}_\Psi \mu(\cdot, r_0) \rangle}{\log r_0} + o\left(\frac{1}{\log r_0}\right) \quad (9)$$

where  $\langle \mathcal{T}_\Psi \mu(\cdot, r_0) \rangle$  is the average value of the wavelet projection over the whole signal, and  $r_0$  is chosen to diminish the relative amplitude of the correction term  $o\left(\frac{1}{\log r_0}\right)$ . When the values of  $h(\vec{x})$  are computed for all points  $\vec{x}$  in the image's domain, they can be sorted and the typical threshold 0 is naturally chosen to define the standard reconstruction set in the MSM method. Such a set often provides a very good reconstruction. Consequently the basic edge set computed automatically in the MSM edge detection method, that yields a good reconstruction, and as such can define the support of a complete compact representation, is defined by:

$$\mathcal{E}_{\text{MSM}} = \bigcup_{h_\infty \leq h \leq 0} \mathcal{F}_h \quad (10)$$

Table 1: **Images used for our experiments.** **Row 1** (from left to right): Harrington weave (SIPI image database), Hibiscus (SIPI image database), car (CMU image database), Lena (SIPI image database), imk01310 (van Hateren image database), imk03324 (van Hateren image database), imk04050 (van Hateren image database). **Row 2:** visualization of the singularity exponents. **Row 3:** compact representation of MSM points corresponding to 30% pixel density.



In general,  $\mathcal{E}_{\text{MSM}}$ , as well as  $\mathcal{F}_{\infty}$  are multifractal, so that a notion of unitary orthogonal field to these sets can be defined only statistically across the scales. A normal unitary field to  $\mathcal{E}_{\text{MSM}}$  is computed in Fourier space by prefiltering it with nonlinear kernel  $\widehat{W}(\vec{k}) = \frac{1}{\|\vec{k}\|}$  corresponding to the convolution with wavelets from the family  $\Psi^{\beta}(\vec{x}) = 1/(1 + |\vec{x}|^2)^{\beta}$ , for  $\beta = 1, 2, 3, 4$ , and averaging the resulting coefficients. Images used for our experiments and their corresponding MSM points are shown in Table 1. For the next part of the paper, we will call our edge detection algorithm corresponding to equations (9) and (10) as  $\mathbf{I}_{\infty}$ .

### 2.3. Reconstructible Systems

In this section, we are led to find mathematically a functional  $\mathcal{G} : L_{\mathbb{C}}^1(\mathbb{R}^2) \rightarrow L_{\mathbb{C}}^1(\mathbb{R}^2)$  which permits the reconstruction of the signal's gradient measure from its restriction to the MSM  $\mathcal{F}_{\infty}$ . The choice of the functional is made under the conditions of verifying the properties of being deterministic, linear, translationally-invariant, isotropic and yield correct power spectrum of natural images [63]. The set  $\mathcal{F}_{\infty}$  being equal to the set of most unpredictable points, the functional  $\mathcal{G}$  is reconstructing the whole gradient  $\nabla I(\vec{x})$  from

the knowledge of  $\nabla I|_{\mathcal{F}_\infty}$  (i.e, the gradient restricted to the set of most unpredictable points). By using the Riesz integral representation theorem for linear continuous functionals and using translational invariance one gets

$$\nabla I(\vec{x}) = \int_{\mathcal{F}_\infty} \nabla I(\vec{y}) G(\vec{x} - \vec{y}) d\vec{y} \quad (11)$$

where  $\int_{\mathcal{F}_\infty} d\vec{y}$  means integration over the MSM and we represent the linear operator  $\mathcal{G}$  by means of its density matrix, denoted by  $G(\vec{x}, \vec{y})$ , as given by the theory of linear representation.

The above equation can be rewritten in the form of a distributional field which contains all the functional (the values of the gradient  $\nabla I$ ) and geometrical (the MSM) information which is specific to the signal to be reconstructed. The essential gradient of the signal  $\nabla_\infty I(\vec{x})$  is defined as the following distribution:

$$\nabla_\infty I(\vec{x}) = \nabla I(\vec{x}) \delta_{\mathcal{F}_\infty}(\vec{x}) \quad (12)$$

where  $\delta_{\mathcal{F}_\infty}(\vec{x})$  is the density of the gradient measure restricted to the MSM. Equation (12) results in a standard convolution as the integration is now performed over all the space and not restricted to the MSM. A practical expression of the reconstruction formula is given by [63]:

$$I(\vec{x}) = \int \langle \vec{g}(\vec{x} - \vec{y}) | \nabla_\infty I(\vec{y}) \rangle d\vec{y} = \vec{g} * \nabla_\infty I(\vec{x}) \quad (13)$$

for a vector  $\vec{g}$ , that is to say the form of a diffusion kernel, but integrated over the multifractal set of most unpredictable points  $\mathcal{F}_\infty$ , and is expressed in the Fourier domain as:

$$\hat{I}(\vec{\omega}) = \langle \hat{\vec{g}}(\vec{\omega}) | \widehat{\nabla_\infty I}(\vec{\omega}) \rangle \quad (14)$$

where  $\langle \cdot | \cdot \rangle$  denoting the standard Hermitian product on  $\mathbb{C}^2$  and  $\widehat{\phantom{x}}$  represents the Fourier transform.

#### 2.4. Determination of the kernel

To derive an exact formulation for the reconstruction kernel  $\hat{\vec{g}}$ , we use an alternate approach to [63] and follow the derivation presented in [64] in the case of differentiable fields. For that matter, given a predefined smooth vector

field  $\vec{v}$ , one can seek an image  $I(\vec{x})$  such that its gradient  $\nabla I(\vec{x})$  minimizes the  $L^2$  distance with  $\vec{v}$  as

$$\underset{I}{\operatorname{argmin}} \int \int \|\nabla I(\vec{x}) - \vec{v}(\vec{x})\|_2^2 d\vec{x} \quad (15)$$

Following the derivation explained in [64] we get a version of the reconstruction kernel  $\hat{g}$  in the smooth case as follows. We take the Euler-Lagrange variational formulation of equation (15) and obtain

$$\operatorname{div}(\nabla I)(\vec{x}) = \operatorname{div}(\vec{v})(\vec{x}) \quad (16)$$

Taking the Fourier transform of equation (16), one gets

$$\hat{I}(\vec{\omega}) = -i \frac{\omega_x \hat{v}_x(\vec{\omega}) + \omega_y \hat{v}_y(\vec{\omega})}{\omega_x^2 + \omega_y^2} \quad (17)$$

where the vector field  $\vec{v}(\vec{x})$ , after Fourier transformation gives rise to a complex vector field  $\hat{\vec{v}} = (\hat{v}_x(\vec{\omega}), \hat{v}_y(\vec{\omega}))$ . Considering equation (14), this suggests the kernel as

$$\hat{g}(\vec{\omega}) = \frac{\vec{\omega}}{i\|\vec{\omega}\|^2} \quad (18)$$

and consequently the following expression of the reconstruction formula in the Fourier domain:

$$\hat{I}(\vec{\omega}) = \frac{\langle \vec{\omega} | \widehat{\nabla_\infty I}(\vec{\omega}) \rangle}{i\|\vec{\omega}\|^2} \quad (19)$$

Fourier inversion of this formula gives the reconstruction of the image from the restriction of the gradient field to the MSM.

### 3. Edge consistency across the scales

Edges are primary features naturally associated to scale invariant properties of natural images, specifically in the case of turbulent signals where symmetry is restored only in statistical sense [24, 31]. As a consequence, the algorithms used in computing edge features should be consistent across the scales: if one has different acquisitions of a same phenomenon at different scales, the resulting edge pixels computed by these algorithms must produce matching edge pixels accross the scales. To check this, we set up an experiment where a same signal at different resolution is generated and

their outputs produced by some classical edge detection algorithms are evaluated. To generate the different resolutions, we use two methods. The first one consists in computing a multiscale version of the signal by using a standard Haar discrete wavelet transform [58]. The second one consists in using the well-known linear scale-space representation developed by Lindeberg *et al* [60, 61]. A linear scale-space family  $L(., t)$  associated to an original signal  $f$  is obtained by convolving  $f$  with a Gaussian kernel  $g$  such that:

$$L(., t) = g(., t) * f \quad (20)$$

where the Gaussian kernel is given by

$$g(x, t) = \frac{1}{2\pi t} e^{-(x^2+y^2)/2t} \quad (21)$$

and  $t > 0$  is the scale parameter.

We take two images: the clock image from SIPI image database [66], which is a standard rigid object, and an excerpt of a remotely sensed image of sea surface temperature (SST) acquired from space. The SST image corresponds to acquisition of a turbulent phenomenon, for which the existence of a multiscale hierarchy comes from the turbulence associated to Navier-Stokes equations [31]. The results for the SST image are shown in Table 2. In this table are displayed (left column) the SST image at various resolutions using the Haar discrete wavelet transform. Columns 2, 3 and 4 show the edge pixels produced by three different algorithms: the algorithm called  $\mathbf{I}_\infty$ , which is the proposed edge detection algorithm (explained in section 2.2), classical Canny edge detector (column 3) and the Mallat-Zhong edge detection technique [38], also known as multiscale Canny edge detection (which is also related to the WTMM method). It is clear from these results that Canny edge pixels are not consistent across the scales: first the boundary of the main coherent structure (the temperature front depicted by the boundary between the dark and light area) is not properly described by Canny edge pixels across the scales. Second, it is difficult to match the corresponding edge pixels across the scales. This can, however, be related to the fact that Canny edge detector encodes all edges regardless of scale, and that is why we also use multiscale Canny in our comparison. Considering the multiscale Canny edge pixels, they behave more consistently compared to Canny edge pixels, but they are still outperformed by  $\mathbf{I}_\infty$ , the latter being specifically designed to retain consistency across the scales. This is confirmed by the results displayed in columns 5, 6, 7 of Table 2, which shows the reconstruction



Table 2: **Inconsistent behaviour of edges along the scales.** Image used is an excerpt from the sea surface temperature (SST) image (MODIS data) of the Agulhas current below the coast of South Africa. Different resolutions of the SST image are formed by taking the approximation coefficients resulting from a standard Haar discrete wavelet transform. The edge pixels are marked red or blue according to the sign of the scalar product between the normal to the set of edge and the image gradient at that point (see section 2.2 after equation (10)). We show the results of the following edge detection algorithms:  $\mathbf{I}_\infty$  (proposed edge detection algorithm, explained in section 2.2), Canny and multiscale Canny (mC) edge detection [38]. D corresponds to the pixel density of the respective edges. For edges computed using Canny edge detector,  $\alpha$  specifies the lower sensitivity threshold and  $\sigma$  is the standard deviation of the Gaussian filter. R- $\mathbf{I}_\infty$ , R-Canny and R-mC are the reconstructions, with equation (19), from their respective edges.


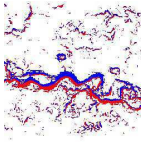
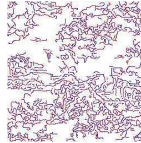
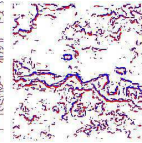
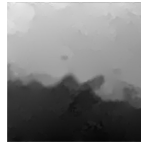
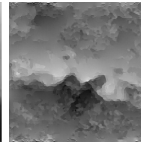
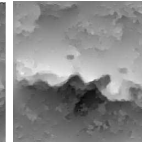

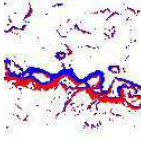
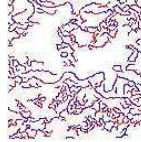
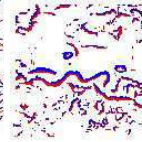

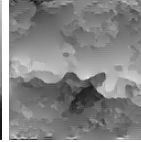
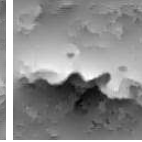

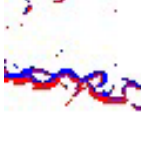
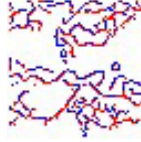
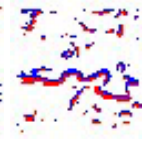

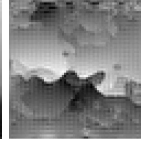
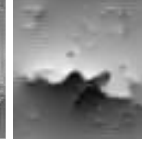

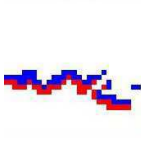

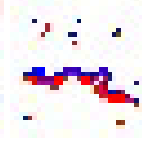
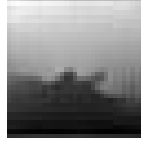
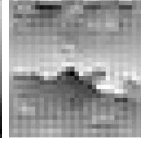
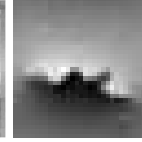

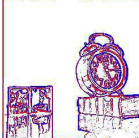
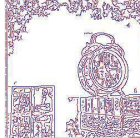
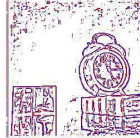






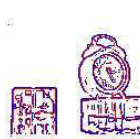




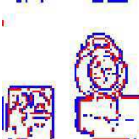





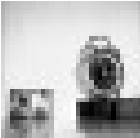
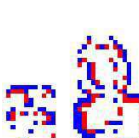


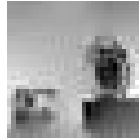
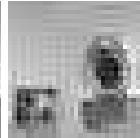
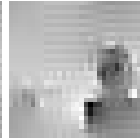
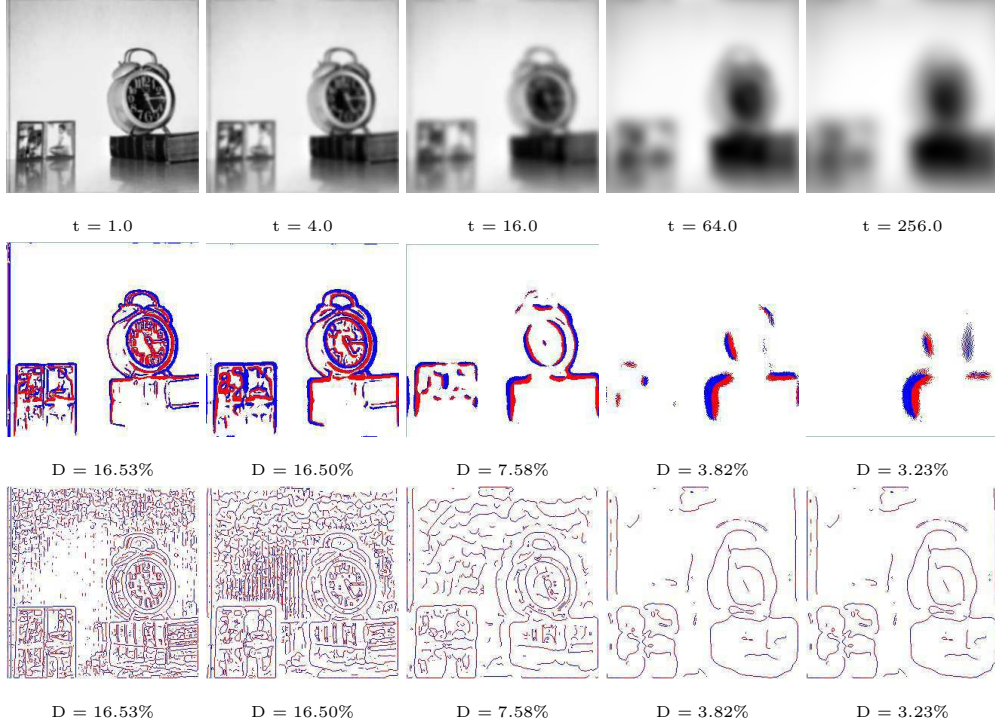
Original	$\mathbf{I}_\infty$	Canny	mC	R- $\mathbf{I}_\infty$	R-Canny	R-mC
						
256 × 256 pixels	D = 16.24%	D = 16.24%	D = 16.98%	psnr=23.46 dB	psnr=4.18 dB	psnr=6.76 dB
		$\alpha=0.1, \sigma=0.03$				
						
128 × 128 pixels	D = 17.45%	D = 17.45%	D = 17.96%	psnr=22.31 dB	psnr=4.50 dB	psnr=6.49 dB
		$\alpha=0.13, \sigma=0.03$				
						
64 × 64 pixels	D = 16.89%	D = 16.89%	D = 17.91%	psnr=20.65 dB	psnr=3.77 dB	psnr=6.69 dB
		$\alpha=0.15, \sigma=0.03$				
						
32 × 32 pixels	D = 18.55%	D = 18.55%	D = 19.04%	psnr=17.81 dB	psnr=3.52 dB	psnr=6.30 dB
		$\alpha=0.16, \sigma=0.03$				

Table 3: **Inconsistent behaviour of edges along the scales.** Image used is the clock image from SIPI image database [66]. Different resolutions of the clock image are formed by taking the approximation coefficients of Haar discrete wavelet transform.

Original	$\mathbf{I}_\infty$	Canny	mC	$\mathbf{R-I}_\infty$	R-Canny	R-mC
						
256 × 256 pixels	D = 17.19%	D = 17.19%	D = 17.35%	psnr= 8.92 dB	psnr= 5.84 dB	psnr= 7.99 dB
$\alpha=0.03, \sigma=0.03$						
						
128 × 128 pixels	D = 16.86%	D = 16.86%	D = 16.62%	psnr= 8.90 dB	psnr= 6.26 dB	psnr= 7.58 dB
$\alpha=0.03, \sigma=0.03$						
						
64 × 64 pixels	D = 16.74%	D = 16.74%	D = 17.09%	psnr= 8.52 dB	psnr= 7.16 dB	psnr= 6.05 dB
$\alpha=0.05, \sigma=0.03$						
						
32 × 32 pixels	D = 16.21%	D = 16.21%	D = 16.40%	psnr= 7.95 dB	psnr= 4.94 dB	psnr= 4.58 dB
$\alpha=0.3, \sigma=0.03$						

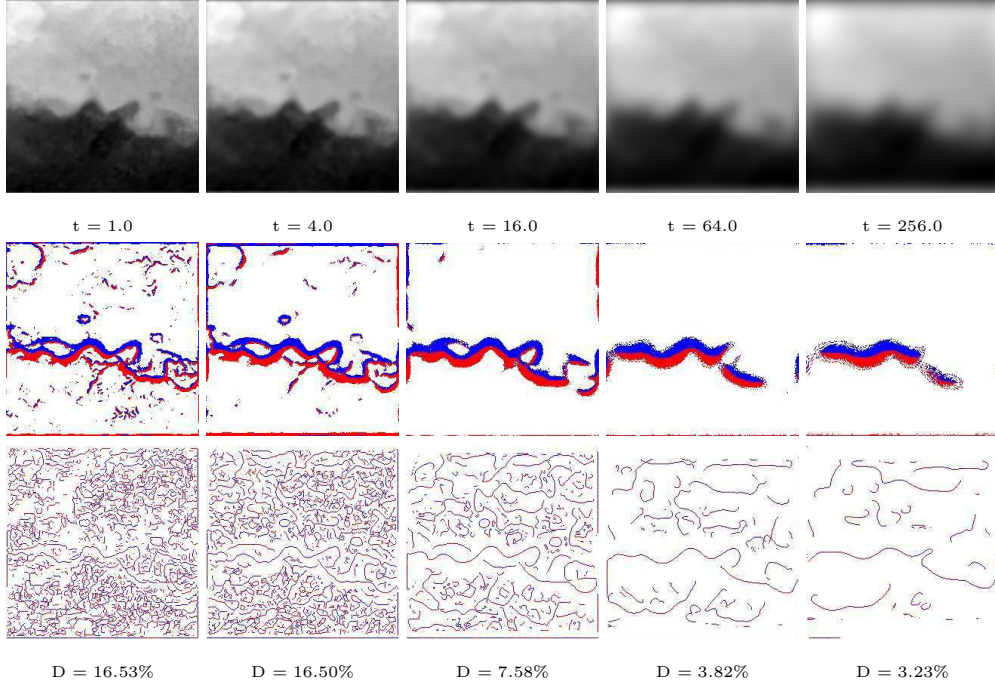
of the whole signal from its gradient measure restricted to the edge pixels. As explained later in this article, this is a crucial test, as edges are supposed to encode in the most efficient way the information content in the signal. The precise reconstruction process shown in this table is fully explained in section 2.3. Note that in Table 2, edge pixels produced by  $\mathbf{I}_\infty$  algorithm encode in a particularly efficient way the main boundary of the temperature front across the scales. The same is reproduced in Table 3 on the clock image,

Table 4: **Evaluation of edge consistency across scales using Lindeberg [60] scale space representation.** **Row 1** (from left to right): different resolutions of the clock image obtained by changing the scale parameter  $t$  (explained in section 3). **Row 2**: compact representation of  $\mathbf{I}_\infty$  points. **Row 3**: edges computed using Lindeberg edge detector.



with similar results. In Table 4 and Table 5, we make a similar comparison between the  $\mathbf{I}_\infty$  edges and the edges produced by the multiscale edge detection of Lindeberg [60]. Here again, we see that in both cases, consistency across the scales is better achieved by  $\mathbf{I}_\infty$  algorithm. In Table 5, the density of the  $\mathbf{I}_\infty$  edges are kept around 16% precisely; but because Lindeberg edges uses a zero-crossing of image intensity further filtered by a third order directional derivative, the number of candidate edge pixels is automatically reduced with increased blurring, resulting in a limited density across the scales. From these experiments we see that, specially in the case of acquired turbulent phenomena, some classical edge detection algorithms, even those based on linear scale-space theory, do not behave consistently across the scales.

Table 5: **Evaluation of edge consistency across scales using Lindeberg scale space representation.** **Row 1** (from left to right): different resolutions of the SST image obtained by changing the scale parameter  $t$ . **Row 2**: compact representation of  $\mathbf{I}_\infty$  points (pixel density fixed at 16%). **Row 3**: edges computed using Lindeberg edge detector.



In this context, one might argue about a comparison with Elder-Zucker's algorithm [49]. Elder-Zucker's algorithm on the detection of edges is inspired by its own local scale control method that determines a unique scale, as a function of filter scale  $\sigma$ , for local estimation at each point in an image [49, 50]. Therefore, a *minimum reliable scale* i.e., a filter with smallest standard deviation  $\sigma$  that can be used reliably, is being determined for each point in the gradient map and used locally to derive logical inferences from derivative estimates of the signal (that are key to edge detection). The choice of  $\sigma$  and the definition of reliability rest with the prior computation of a *critical value function* that depends on some statistical parameters [57]. A good choice of these statistical parameters are necessary for the proper functioning of the algorithm; incorrect estimate may lead to reduced performances. The performance of  $\mathbf{I}_\infty$ , on the other hand, is independent of the choice of any  $a$

*priori* inputs and is applied directly on different scale-based representation of images (without doing any processing to enhance the performance of edge detection) and verify its performance across scales. We are interested in extracting the most singular components that are related to edges in an image (by proper computation of singularity exponents on the image, as explained in section 2.2), irrespective of the scale or the spatial representation of the image.

## 4. Results

### 4.1. Choice of images

Images of different entities are chosen trying to cover a broad spectrum of natural images starting from textures (Harrington weave) to an object (car), face (lena) to landscapes (imk01310, imk03324) and signals of turbulent acquisitions (sea surface temperature, turbulence degraded atmospheric phase) as is shown in Table 1 and Table 7. These experimental images are chosen from standard databases like SIPI image database (Harrington weave, Hibiscus, Lena) [66], CMU image database (car) [67] and the van Hateren database (imk01310, imk03324, imk04050) [27]. The SST data was acquired by MODIS satellite and the perturbed phase data is provided by the french aerospace lab ONERA.

### 4.2. Experimental procedure

We make use of natural images described in section 4.1 for our experiments. The classical edge algorithms tested are Matlab<sup>®</sup> implementations. Reconstruction is performed from edges, calculated using different edge detector algorithms, using equation (19). The pixel densities of all the edges calculated using different edge detectors are kept within a close range (given the constraint imposed by the Gaussian sigma), with the pixel density of  $\mathbf{I}_\infty$  kept lower than the others. Similar experiment is repeated over the images by adding different proportions of Gaussian white noise with respect to the standard deviation of the input images (see Table 8 and Table 10).

### 4.3. Quantitative analysis of the results

The images we use, together with their reconstruction using different edge detectors, are shown in Table 6 and Table 7. Performance of the reconstruction under different levels of noise (SNR = 26 dB, 14 dB and 6 dB) is shown in Table 8. Visual quality of the reconstructed images, shown in Table 6,



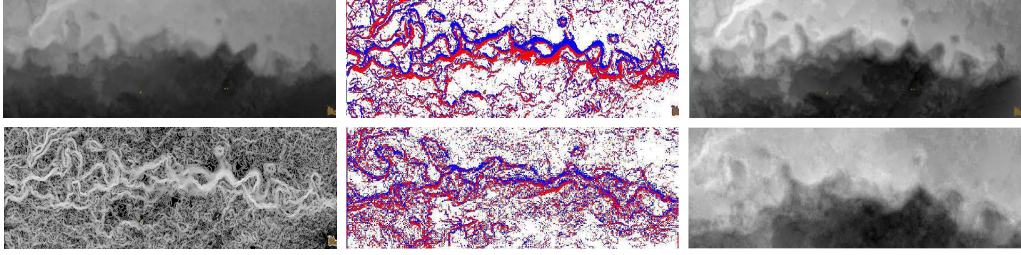


Figure 1: This figure illustrates the complexity of edge detection in the case of turbulent images. **Top row** (from left to right): excerpt of the Agulhas current from the SST image, set of edge pixels computed by  $\mathbf{I}_\infty$  corresponding to 25% of pixel density in the selected area, reconstruction with algorithm  $\mathbf{I}_\infty$ . **Bottom row** (from left to right): singularity exponents of the SST image, edges produced by algorithm NLFS [53] which behaves the best among the classical edge operators tested, reconstruction with edge pixels produced by NLFS. The coherent structures are not respected by NLFS, showing the superiority of  $\mathbf{I}_\infty$ .

Table 7 and Table 8, are evaluated based on the structural similarity index metric (SSIM) [68]. SSIM's for the reconstructed images shows that  $\mathbf{I}_\infty$  outperforms the classical edge detectors in majority of the cases.

We also illustrate the pertinence of the framework of reconstructible systems for evaluating an edge operator's compact representation effectiveness using the *mean square error* (MSE) and *peak signal to noise ratio* (PSNR, expressed in decibels dB) defined by:

$$\text{MSE} = \frac{1}{m \times n} \sum_{i,j} |I(x_{i,j}) - I_r(x_{i,j})|^2 \quad (22)$$

$$\text{PSNR} = 20.0 \times \log_{10} \frac{\max(I(\vec{x}))}{\sqrt{\text{MSE}}} \quad (23)$$

Table 9 shows the performance evaluation of the reconstructed images using PSNR and MSE metric and Table 10 shows the performance of the reconstruction under different levels of input SNR using the same metrics.

#### 4.4. Comments: The case of turbulent signals

Examination of the results for SST and phase (turbulent phenomena) images are particularly interesting: an edge in a turbulent signal is poorly

characterized by a filter’s response to step functions, and the case of Fully Developed Turbulence is paradigmatic for the existence of a multiscale hierarchy associated to cascading dynamics of physical variables [31]. Tuning with the scale-space parameter given by the kernel’s standard deviation modifies the input signal incoherently w.r.t. to the real multiscale hierarchy present in the data, resulting in bad reconstructions for most Gaussian prefiltered edge operators, and with no enhancing by adding more edge pixels. Universality classes are not well characterized by the transitions associated to classical edge operators. Among these operators, NLFS seems to behave the best, yet  $\mathbf{I}_\infty$  outperforms it, and provides also an automatic threshold determination for the set of edges. In Fig 1, we show an extended part of the SST data and illustrate the sensitivity of edge detection and signal reconstruction between  $\mathbf{I}_\infty$  and NLFS.

## 5. Discussion

In this article, we presented a performance evaluation of the quality of edge pixels based on the reconstruction formula presented in section 2.3. It happens that  $\mathbf{I}_\infty$  edge pixels lead, for most images used in this work, the best quantitative results in terms of PSNR, SSIM and MSE. But  $\mathbf{I}_\infty$  edge pixels are defined irrespective of any reconstruction formula. Consequently, the evaluation procedure based on the reconstruction formula is independent of the definition of edge pixels. In fact, it is easy to see that any reconstruction formula designed to test the accuracy of a compact representation must have, under reasonable hypothesis the form displayed in equation (19). Indeed, if one suppose that a reconstructor is able to generate the whole signal from its set  $\mathcal{K}$  of edges, using a linear functional  $\mathcal{G}$ , then equation (11) must be valid, with  $\mathcal{F}_\infty$  replaced by  $\mathcal{K}$ . Then assuming linearity, translational invariance and isotropy, one gets the following reconstruction formula for testing the validity of the set  $\mathcal{K}$  of edge pixels:

$$\widehat{I(\vec{\omega})} = \frac{\langle \vec{\omega} | \widehat{\nabla I|_{\mathcal{K}}} \rangle}{i \|\vec{\omega}\|^2} \quad (24)$$

In that sense, we can say that the reconstruction technique presented in this work is universal and not specifically designed for the assumptions motivating the article and suitable for the proposed approach.

## 6. Conclusion

Tested classical edge detection algorithms are reviewed in this work in the sense that they lack compact representation. When considering images as acquisitions of a complex system with unknown extended phase space, edge pixels correspond to high-order transitions, which are characterized by critical exponents. These exponents define universality classes, they describe common statistical and macroscopic multiscale characteristics of the system. When they are numerically evaluated at high precision through wavelet projection measures in a microcanonical setting, they allow the determination of the multiscale hierarchy, of multifractal nature, which is very different to the scale-space models commonly used. The framework of reconstructible system permits a quantitative evaluation of the compactness of a representation, leading to extremely accurate edge pixels, specifically in the case of turbulent data, where classical edge detectors fail. The introduced edge operator  $\mathbf{I}_\infty$  outperforms on natural images and geophysical variables.

## Acknowledgment

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Table 6: **Results of reconstruction over different edge detection algorithms.** Image description available in Table 1.

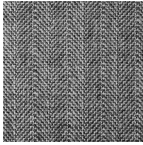
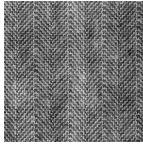
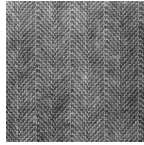
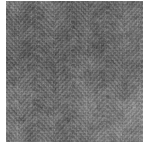
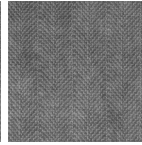
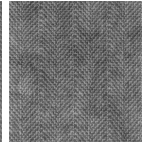
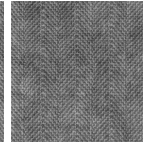



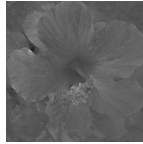
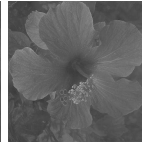
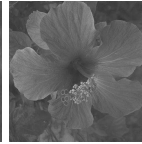
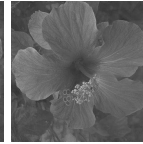
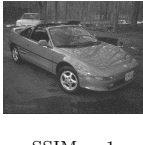
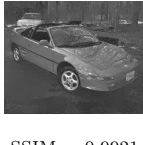



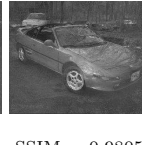
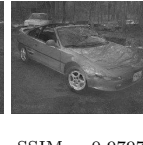

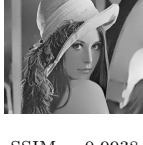
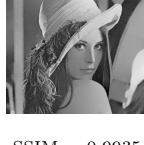
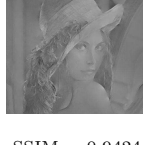
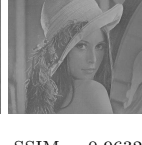
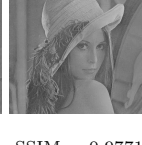


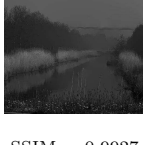
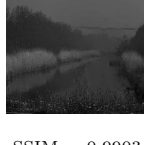
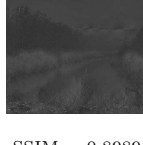
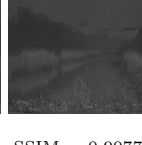
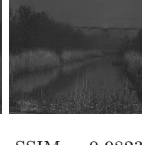

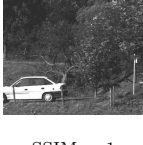

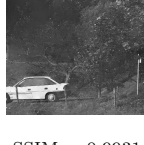





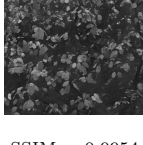
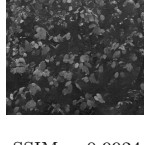
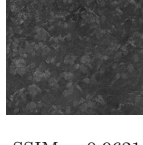
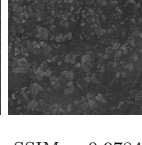
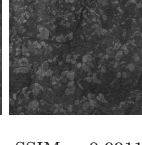

Original	$I_\infty$	NLFS	Canny	LoG	Sobel	Prewitt
						
SSIM = 1	SSIM = 0.9881	SSIM = 0.9848	SSIM = 0.9464	SSIM = 0.9839	SSIM = 0.9837	SSIM = 0.9844
						
SSIM = 1	SSIM = 0.9899	SSIM = 0.9886	SSIM = 0.9683	SSIM = 0.9773	SSIM = 0.9862	SSIM = 0.9863
						
SSIM = 1	SSIM = 0.9921	SSIM = 0.9879	SSIM = 0.9347	SSIM = 0.9624	SSIM = 0.9805	SSIM = 0.9797
						
SSIM = 1	SSIM = 0.9938	SSIM = 0.9935	SSIM = 0.9424	SSIM = 0.9632	SSIM = 0.9771	SSIM = 0.9769
						
SSIM = 1	SSIM = 0.9927	SSIM = 0.9903	SSIM = 0.8989	SSIM = 0.9077	SSIM = 0.9823	SSIM = 0.9848
						
SSIM = 1	SSIM = 0.9925	SSIM = 0.9931	SSIM = 0.9443	SSIM = 0.9773	SSIM = 0.9879	SSIM = 0.9873
						
SSIM = 1	SSIM = 0.9954	SSIM = 0.9924	SSIM = 0.9621	SSIM = 0.9784	SSIM = 0.9911	SSIM = 0.9913

Table 7: **Reconstruction results from different edge detection algorithms over turbulent data.**





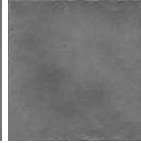
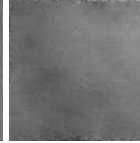
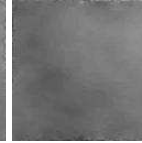


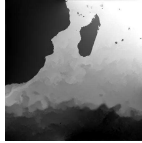
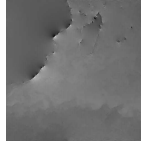
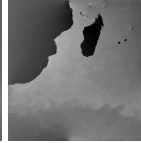
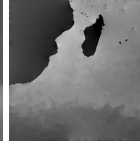
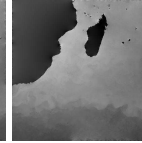
Original	$I_\infty$	NLFS	Canny	LoG	Sobel	Prewitt
						
SSIM = 1	SSIM=0.9986	SSIM=0.9983	SSIM=0.9293	SSIM=0.9026	SSIM=0.9886	SSIM=0.9878
						
SSIM = 1	SSIM=0.9988	SSIM=0.9957	SSIM=0.9182	SSIM=0.9202	SSIM=0.9665	SSIM=0.9881

Table 8: **Performance under noise:** reconstruction results on different edge detection algorithms under different levels of Gaussian white noise.



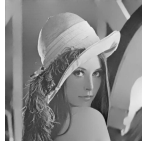




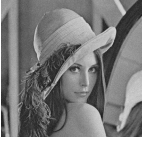
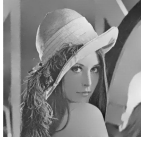
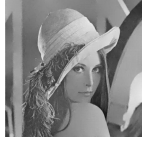
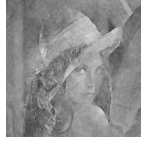




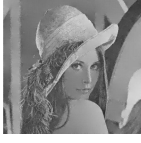





Original	$I_\infty$	NLFS	Canny	LoG	Sobel	Prewitt
						
SNR = 26 dB	psnr=20.94 dB	psnr=20.74 dB	psnr=10.75 dB	psnr=12.83 dB	psnr=14.88 dB	psnr=14.69 dB
	SSIM=0.9937	SSIM=0.9934	SSIM=0.9407	SSIM=0.9616	SSIM=0.9764	SSIM=0.9758
						
SNR = 14 dB	psnr=18.42 dB	psnr=18.00 dB	psnr=10.60 dB	psnr=12.40 dB	psnr=13.71 dB	psnr=14.21 dB
	SSIM=0.9932	SSIM=0.9927	SSIM=0.9343	SSIM=0.9611	SSIM=0.9722	SSIM=0.9755
						
SNR = 06 dB	psnr=15.82 dB	psnr=15.76 dB	psnr=10.02 dB	psnr=12.18 dB	psnr=13.26 dB	psnr=13.47 dB
	SSIM=0.9895	SSIM=0.9886	SSIM=0.9206	SSIM=0.9524	SSIM=0.9650	SSIM=0.9686

Table 9: **Quantitative analysis:** evaluation of the reconstruction over different edge detection algorithms, in terms of PSNR and MSE, corresponding to Table 6 and Table 7. The pixel density (D) of all the edges are kept within a close range, with the pixel density of  $\mathbf{I}_\infty$  being the lowest.

Image	Params	$\mathbf{I}_\infty$	NLFS [53]	Canny	Lindeberg [60]	LoG	Sobel	Prewitt
weave (1024 × 1024 pixels)	D (%) mse psnr (dB)	31.5 0.1439 <b>17.38</b>	32.87 0.1977 16.01	29.56 0.6600 10.76	31.86 0.3955 13.99	31.28 0.2790 14.51	31.57 0.2148 15.64	31.67 0.2158 15.62
hibiscus (512 × 512 pixels)	D (%) mse psnr (dB)	28.55 0.0906 <b>22.43</b>	28.87 0.0944 21.92	28.55 0.4005 16.04	29.15 0.3139 17.04	29.98 0.2657 17.76	29.43 0.1690 19.72	29.22 0.1622 19.80
car (260 × 320 pixels)	D (%) mse psnr (dB)	30.00 0.0794 <b>25.59</b>	36.03 0.1321 23.37	32.28 0.7350 15.84	36.28 0.3972 18.68	37.58 0.4038 18.53	36.86 0.2150 21.26	37.24 0.2262 21.04
Lena (512 × 512 pixels)	D (%) mse psnr (dB)	32.00 0.0563 <b>21.17</b>	32.93 0.0587 20.99	28.80 0.5898 10.98	32.20 0.3718 12.98	32.36 0.3775 12.91	30.13 0.2298 15.06	30.26 0.2305 15.05
imk01310 (512 × 512 pixels)	D (%) mse psnr (dB)	30 0.0626 <b>29.43</b>	30.12 0.0923 27.74	27.98 1.1389 16.88	31.09 1.0744 17.08	31.24 0.9434 17.65	29.90 0.2150 24.07	30.40 0.1933 24.53
imk03324 (512 × 512 pixels)	D (%) mse psnr (dB)	30.00 0.0854 <b>23.08</b>	30.14 0.1057 22.15	32.48 0.7808 13.42	35.43 0.3879 16.51	36.46 0.3168 17.38	36.06 0.1634 20.26	36.82 0.1705 20.07
imk04050 (512 × 512 pixels)	D (%) mse psnr (dB)	31.00 0.0652 <b>25.79</b>	31.32 0.0974 24.05	31.22 0.4634 17.23	32.59 0.3554 19.17	32.95 0.2996 22.60	32.61 0.1358 20.50	32.93 0.1361 22.59
phase (128 × 128 pixels)	D (%) mse psnr (dB)	25.00 0.0184 <b>23.97</b>	27.27 0.0187 23.14	24.63 0.8025 7.08	23.62 1.3023 5.29	29.15 0.9307 5.34	25.91 0.1660 14.24	25.42 0.1745 14.02
SST (512 × 512 pixels)	D (%) mse psnr (dB)	25.00 0.0114 <b>23.99</b>	30.73 0.0404 18.50	21.62 0.8105 5.56	22.64 0.8257 5.40	23.95 0.7545 5.79	23.80 0.2035 11.48	24.01 0.2011 11.54

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Table 10: **Quantitative analysis (noisy environment)**: evaluation of the reconstruction over different edge detection algorithms, in terms of PSNR and MSE, under different levels of noise. For all images, the first, second and third row represents the performance under input SNR of 26 dB, 14 dB and 6 dB respectively.

Image	$I_{\infty}$	NLFS	Canny	Lindeberg	LoG	Sobel	Prewitt
	psnr , mse	psnr , mse	psnr , mse	psnr , mse	psnr , mse	psnr , mse	psnr , mse
weave	16.89 , 0.1763	15.49 , 0.2384	10.37 , 0.7858	13.34 , 0.4176	14.45 , 0.2964	15.60 , 0.2278	15.51 , 0.2356
	16.68 , 0.2388	15.21 , 0.3301	10.05 , 0.8449	13.12 , 0.4986	14.24 , 0.4103	15.48 , 0.3344	15.22 , 0.3245
	14.51 , 0.5388	14.61 , 0.6217	9.75 , 1.0993	12.78 , 0.8647	13.53 , 0.7520	13.98 , 0.6500	13.57 , 0.6169
hibiscus	22.18 , 0.0954	21.88 , 0.1078	16.01 , 0.4336	16.76 , 0.3277	17.72 , 0.2755	19.69 , 0.1762	19.66 , 0.1728
	20.64 , 0.1553	20.24 , 0.1708	15.26 , 0.5146	16.11 , 0.4350	16.93 , 0.3385	18.96 , 0.2590	18.29 , 0.2685
	17.66 , 0.4300	16.07 , 0.4420	13.49 , 0.9930	14.81 , 0.6040	14.20 , 0.8359	15.53 , 0.6033	16.24 , 0.6487
car	25.48 , 0.0837	22.87 , 0.1535	15.80 , 0.7351	18.57 , 0.4000	18.52 , 0.4040	21.22 , 0.2260	20.99 , 0.2364
	23.69 , 0.1416	22.05 , 0.2189	15.72 , 0.7951	18.11 , 0.4668	18.12 , 0.4612	20.81 , 0.2796	20.02 , 0.3071
	19.41 , 0.3952	17.96 , 0.4794	15.40 , 1.0349	17.67 , 0.5550	16.61 , 0.7574	18.39 , 0.9708	17.98 , 0.5778
imk01310	29.23 , 0.0627	27.53 , 0.0951	16.82 , 1.1749	16.96 , 1.0853	17.57 , 0.9494	23.83 , 0.2232	24.46 , 0.1938
	26.21 , 0.1247	25.23 , 0.1595	15.98 , 1.3719	16.18 , 1.1278	17.45 , 0.9709	22.34 , 0.3155	22.16 , 0.2756
	19.64 , 0.3872	18.75 , 0.4196	13.28 , 1.5785	15.33 , 1.2092	14.23 , 1.3095	16.67 , 0.6838	17.99 , 0.6100
imk03324	23.07 , 0.0921	21.98 , 0.1216	13.39 , 0.8802	16.32 , 0.3941	17.34 , 0.3571	20.10 , 0.1932	19.43 , 0.2217
	22.16 , 0.1551	21.80 , 0.1568	13.13 , 0.8853	16.08 , 1.1278	17.15 , 0.4636	19.52 , 0.2753	19.33 , 0.2911
	18.82 , 0.4310	18.22 , 0.4578	12.10 , 0.9979	15.80 , 0.5712	16.13 , 0.6700	18.13 , 0.5482	17.40 , 0.5882
imk04050	25.53 , 0.0696	23.95 , 0.1005	17.21 , 0.4712	18.35 , 0.3591	18.95 , 0.3149	22.49 , 0.1429	22.45 , 0.1422
	23.25 , 0.1317	21.80 , 0.1697	16.46 , 0.5960	17.00 , 0.4874	18.05 , 0.3875	20.60 , 0.2379	19.74 , 0.2580
	18.06 , 0.4295	17.66 , 0.4450	14.35 , 0.9353	14.89 , 0.8137	15.36 , 0.7492	16.59 , 0.6069	16.07 , 0.6592
phase	23.30 , 0.0230	22.77 , 0.0261	7.69 , 0.8889	5.15 , 1.4985	6.55 , 1.0722	13.07 , 0.2603	12.63 , 0.2604
	19.02 , 0.0883	19.35 , 0.0884	7.17 , 0.9437	5.01 , 1.6293	6.00 , 1.2843	10.65 , 0.5678	10.97 , 0.5775
	15.39 , 0.3784	15.87 , 0.3693	6.59 , 0.9051	4.59 , 2.2028	5.03 , 1.3440	9.80 , 1.1675	9.26 , 0.9469
SST	23.47 , 0.0159	18.02 , 0.0427	5.22 , 0.8383	5.06 , 0.8354	5.41 , 0.7707	10.11 , 0.3280	10.74 , 0.2852
	19.21 , 0.0822	17.85 , 0.0979	5.09 , 0.9554	4.88 , 0.8536	4.96 , 0.8102	9.74 , 0.4179	10.53 , 0.3253
	15.09 , 0.3686	15.03 , 0.3797	4.86 , 1.3826	4.10 , 0.9447	4.14 , 0.9676	9.00 , 0.7962	9.76 , 0.7958

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